

## The Crystal Structure of Adenosine\*

BY T. F. LAI† AND RICHARD E. MARSH

Arthur Amos Noyes Laboratory of Chemical Physics, California Institute of Technology,  
Pasadena, California 91109, U.S.A.

(Received 9 December 1971)

The crystal structure of adenosine,  $C_{10}H_{13}N_5O_4$ , has been determined from the intensities of 1333 reflections, each measured at least four times on an automated diffractometer. The crystals are monoclinic, space group  $P2_1$ , with  $a = 4.825$  (1),  $b = 10.282$  (2),  $c = 11.823$  (1) Å,  $\beta = 99.30$  (1)°, and two molecules per cell. Least-squares refinement of coordinates and anisotropic temperature factors for all atoms, including hydrogen, led to an  $R$  index of 0.024 and standard deviations of about 0.003 Å for bond distances between pairs of heavy atoms. The anomalous scattering by nitrogen and oxygen was used to confirm the absolute configuration of the sugar ring. The conformation of the ribose ring is  $C(3')$  *endo*; the torsion angle about the glycosidic bond is 9.9°. Intermolecular interactions include a full complement of hydrogen bonds, a relatively short  $C(2)-H \cdots O(2)$  contact (3.09 Å), and parallel stacking of adenine rings with an interplanar spacing of 3.57 Å.

### Introduction

The determination of the crystal structure of adenosine was undertaken as part of a research program on the structures of compounds related to nucleic acids, and with a particular aim towards furnishing accurate details of the molecular geometry and obtaining further information concerning intermolecular forces such as hydrogen bonding and base stacking.

### Experimental

Crystals of adenosine, in the form of colorless prisms, were kindly supplied by Dr C. E. Bugg. Preliminary oscillation and Weissenberg photographs showed them to be monoclinic. The absence of reflections  $0k0$  with  $k$  odd indicated the space group  $P2_1$ ,  $P2_1/m$  being ruled out since the nucleoside is optically active.

Accurate cell dimensions were obtained from a least-squares fit to 93 measurements made on  $h0l$  and  $0kl$  Weissenberg photographs prepared in a Straumanis-type camera. The density was measured by

flotation. Crystallographic data are given in Table 1.

A cubic crystal of dimensions  $0.1 \times 0.1 \times 0.1$  mm was selected for intensity data collection and was mounted in an arbitrary orientation on a Datex-automated General Electric XRD-5 diffractometer. Intensities were measured using Ni-filtered copper radiation and  $\theta-2\theta$  scans at a scanning speed of 1° (in  $2\theta$ ) per min; the background was counted for 30 sec at each extremum. The scan ranges were adjusted to account for  $\alpha_1-\alpha_2$  splitting. All reflections within one quadrant of the sphere of reflection extending to  $2\theta = 155^\circ$  were measured four times; intensities of about 100 very strong reflections were measured a fifth time at a reduced current setting. A check reflection was monitored every twenty reflections; it showed no significant drift in intensity during nearly six weeks of data collection.

Of 5500 measurements, 56 were rejected because they differed distinctly from the other measurements in the set. These rejected measurements usually occurred in groups of two or more, leading us to suspect failures in the instrumentation. The remaining measurements were averaged, with appropriate weights, to yield 1333 intensities, of which 5 were less than one e.s.d. above background. Individual variances  $\sigma^2(I_i)$  were calculated from counting statistics with an additional term  $(0.02 S)^2$ , where  $S$  is the scan count; variances in the averaged intensities,  $\sigma^2(I)$ , were calculated from the expression

$$\sigma^{-2}(I) = \sum_i \sigma^{-2}(I_i).$$

Finally, the intensities and their standard deviations were placed on an absolute scale by Wilson's (1942) method. No correction was made for absorption ( $\mu t \approx 0.1$ ). All calculations were carried out on an IBM 7094 computer under the *CRYRM* system (Duchamp, 1964).

Table 1. *Crystal data*

Adenosine	$C_{10}H_{13}N_5O_4$
Space group: $P2_1$	$F(000) = 280$
$a = 4.825 \pm 0.001$ Å	F.W. 267.25
$b = 10.282 \pm 0.002$	$Z = 2$
$c = 11.823 \pm 0.001$	$D_x = 1.533 \text{ g.cm.}^{-3}$
$\beta = 99.30 \pm 0.01^\circ$	$D_m = 1.54 \text{ g.cm.}^{-3}$
$(\lambda \text{Cu } K\alpha = 1.5418 \text{ \AA})$	

\* Contribution No. 4391 from the Arthur Amos Noyes Laboratory of Chemical Physics. This investigation was supported in part by Public Health Service Research Grant No. GM 16966, from the National Institute of General Medical Sciences.

† Present address: Department of Chemistry, University of Hong Kong, Hong Kong.

### Determination and refinement of the structure

The orientation of the planar adenine moiety was derived from a three-dimensional sharpened Patterson synthesis and approximate  $y$  and  $z$  coordinates for these atoms were deduced from packing and hydrogen bonding considerations. The ribose ring was located on subsequent electron density projections onto (100). Further packing studies and three-dimensional electron density maps then led to a set of approximate  $x$  coordinates. The hydrogen atoms were located from difference maps calculated during the course of the refinement.

Refinement was carried out mainly by least-squares methods. The quantity minimized was  $\sum w(F_o^2 - F_c^2/k^2)^2$ , with weights  $w$  equal to  $1/\sigma^2(F_o^2)$  and  $F_c^2$  as defined by Larson (1967), equation (3). Throughout the refinement the  $y$  coordinate of O(1') was held fixed. The full matrix of normal equations was collected until it became too large (over 207 parameters) to fit into core memory. Atomic scattering factors for the non-hydrogen atoms were obtained from *International Tables for X-ray Crystallography* (1962), and those for hydrogen atoms were from Stewart, Davidson & Simpson (1965).

Eventually, a list of 289 parameters was adjusted: coordinates and anisotropic temperature parameters of all 32 atoms (including hydrogen), a scale factor, and a secondary extinction parameter  $g$  (Zachariasen, 1963; Larson, 1967). The anisotropic temperature factors for the heavy atoms, the scale factor and the secondary extinction parameter were kept in one matrix and the positional parameters in a second matrix; the anisotropic temperature parameters of the hydrogen atoms were adjusted during two least-squares

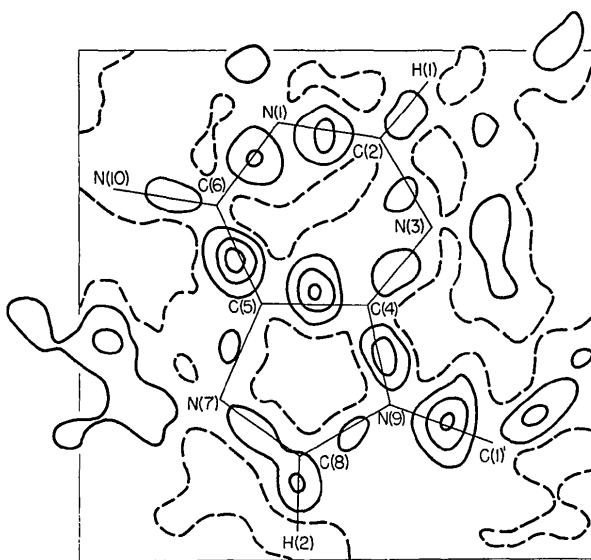


Fig. 1. A portion of the final difference map, evaluated in the plane of the adenine moiety. Contours are at intervals of  $0.05 \text{ e.}\text{\AA}^{-3}$ ; the zero contour is dashed.

cycles in which the remaining parameters were held fixed. The refinement approached convergence at an  $R$  value of 0.025.

In view of the low  $R$  value attained it was thought that the accuracy of the intensity data must be sufficiently high that the anomalous scattering by oxygen and nitrogen could be used to check the absolute configuration of the ribose moiety. Accordingly, structure factors for both enantiomers were calculated, using the following correction terms (Hope, 1969):\*

Atom	$\Delta f'$	$\Delta f''$
N	0.032	0.019
O	0.049	0.032

The  $R$  index was 0.0261 for the correct configuration and 0.0268 for the opposite enantiomer. (No adjustment for the change in scale factor  $k$  was made in these calculations, which accounts for the worsening of  $R$  for both alternatives.)

Refinement was concluded with two least-squares cycles based on the correct enantiomer and with the form factors of nitrogen and oxygen corrected for anomalous dispersion. Coordinates of all the atoms were included in one full matrix while the anisotropic temperature coefficients of the heavy atoms and the scale and extinction parameters were included in a second matrix; the anisotropic coefficients of the hydrogen atoms were not further adjusted. During the last cycles, two reflections 2,12,2 and 303 with residuals  $|F_o^2 - F_c^2|$  over 15 times as large as the standard deviations  $\sigma(F_o^2)$  were given zero weight. For the 2,12,2 reflection, which lies at the outer edge of the effective sphere of reflection, the second background count was appreciably higher than the first, suggesting that the entire  $\alpha_1$ - $\alpha_2$  doublet was not captured; we have no explanation (other than the standard apologia of double reflection) for the discrepancy in the 303 reflection. In the final cycle, no heavy-atom parameter shifted by as much as one-eighth of its e.s.d., nor any hydrogen parameter by 0.3 e.s.d. The final  $R$  index for 1333 independent reflections is 0.024 and the goodness-of-fit is 3.47. The final value of the extinction parameter  $g$  was  $(5.6 \pm 0.2) \times 10^{-5}$ .

At the conclusion of the refinement a three-dimensional difference map was calculated. A section passing through the best plane of the adenine moiety is shown in Fig. 1. Apart from regions between two bonded atoms, no peaks or troughs exceeded  $0.10 \text{ e.}\text{\AA}^{-3}$  in magnitude.

Observed and calculated structure factors are listed in Table 2. Atomic parameters and their standard deviations are listed in Table 3. Since the parameters were partitioned into three separate matrices, these e.s.d.'s may be slightly underestimated. We make no claim that the e.s.d.'s of the anisotropic temperature

\* Slightly different values have more recently been reported by Cromer & Liberman (1970).

THE CRYSTAL STRUCTURE OF ADENOSINE

Table 2. Observed and calculated structure factors ( $\times 10^3$ ) and phase angles ( $^\circ$ )

Table with multiple columns for observed and calculated structure factors and phase angles. Columns include indices (h, k, l), observed values, calculated values, and phase angles in degrees. The table is organized into several groups of reflections.

coefficients of the hydrogen atoms are realistic. These parameters were not adjusted after the anomalous dispersion corrections were made; moreover, a few of them shifted by more than one e.s.d. during the final cycle in which they were adjusted, and it appeared as though convergence would be slow. Nevertheless, we are convinced that these parameters are of some significance. During the two least-squares cycles in which they were adjusted the goodness-of-fit dropped by about 15% below the value obtained with converged isotropic temperature parameters for the hydrogen

atoms. In addition, the resulting thermal ellipsoids are quite reasonable.

The relatively large value of the goodness-of-fit deserves some comment. The individual values of the weighted residuals,  $\sqrt{w(F_o^2 - F_c^2)}$ , appear to be normally distributed, and independent of the magnitude of  $F_o^2$ . However, a disproportionate number of the larger residuals represent low-angle reflections; taken in conjunction with the final difference map (Fig. 1), this trend suggests that the assumption of spherically-symmetric atomic form factors may be an important

Table 3. Final parameters and their standard deviations (in parentheses)

The coordinates and temperature coefficients of the heavy atoms have been multiplied by  $10^4$ ; those of the hydrogen atoms by  $10^3$ . The anisotropic temperature factor is expressed in the form

$$\exp [-(b_{11}h^2 + b_{22}k^2 + b_{33}l^2 + b_{12}hk + b_{13}hl + b_{23}kl)].$$

	<i>x</i>	<i>y</i>	<i>z</i>	<i>b</i> <sub>11</sub>	<i>b</i> <sub>22</sub>	<i>b</i> <sub>33</sub>	<i>b</i> <sub>12</sub>	<i>b</i> <sub>13</sub>	<i>b</i> <sub>23</sub>
N(1)	9236 (4)	-899 (2)	1862 (1)	484 (8)	49 (1)	53 (1)	109 (5)	140 (5)	12 (2)
C(2)	7928 (5)	-722 (2)	2767 (2)	625 (12)	54 (2)	56 (1)	139 (7)	198 (7)	39 (3)
N(3)	6384 (4)	282 (2)	3007 (1)	471 (8)	48 (1)	45 (1)	76 (5)	148 (5)	21 (2)
C(4)	6230 (3)	1189 (2)	2175 (1)	261 (7)	36 (1)	36 (1)	6 (5)	46 (4)	-11 (2)
C(5)	7504 (3)	1165 (2)	1210 (1)	261 (7)	40 (1)	33 (1)	-3 (5)	49 (4)	-2 (2)
C(6)	9096 (3)	41 (2)	1057 (1)	274 (7)	40 (1)	39 (1)	8 (5)	50 (5)	-11 (2)
N(7)	6857 (3)	2275 (2)	557 (1)	380 (8)	51 (1)	44 (1)	45 (5)	99 (4)	16 (2)
C(8)	5226 (4)	2930 (2)	1135 (1)	400 (9)	48 (1)	43 (1)	74 (6)	99 (5)	18 (2)
N(9)	4768 (3)	2340 (2)	2120 (1)	275 (6)	38 (1)	36 (1)	27 (4)	54 (4)	0 (2)
N(10)	10458 (4)	-147 (2)	176 (1)	407 (8)	57 (1)	46 (1)	59 (5)	116 (5)	1 (2)
C(1')	3140 (3)	2782 (2)	2993 (1)	254 (7)	36 (1)	40 (1)	16 (5)	72 (5)	-4 (2)
C(2')	5049 (4)	3161 (2)	4106 (1)	310 (8)	42 (1)	39 (1)	54 (5)	37 (5)	-5 (2)
C(3')	5437 (4)	4616 (2)	3921 (1)	252 (7)	44 (1)	39 (1)	7 (5)	51 (4)	-9 (2)
C(4')	2585 (3)	5024 (2)	3273 (1)	228 (7)	39 (1)	43 (1)	7 (5)	54 (4)	-8 (2)
C(5')	2661 (4)	6177 (2)	2490 (2)	340 (8)	46 (2)	57 (1)	15 (6)	35 (5)	17 (2)
O(1')	1583 (2)	3897	2590 (1)	263 (5)	42 (1)	50 (1)	39 (4)	8 (3)	-12 (2)
O(2')	3579 (3)	3010 (2)	5037 (1)	668 (9)	44 (1)	42 (1)	17 (5)	152 (5)	0 (2)
O(3')	6341 (3)	5337 (2)	4936 (1)	355 (6)	51 (1)	46 (1)	-46 (4)	37 (4)	-21 (2)
O(5')	-83 (3)	6422 (2)	1890 (1)	411 (7)	48 (1)	59 (1)	88 (4)	1 (4)	10 (2)
H(1)	826 (5)	-136 (3)	336 (2)	62 (12)	13 (3)*	10 (2)	0 (11)*	27 (8)	2 (4)*
H(2)	441 (5)	376 (2)	92 (2)	60 (11)	8 (3)	5 (2)	13 (10)	6 (8)	2 (4)
H(3)	1128 (4)	-93 (2)	8 (2)	55 (12)	6 (3)	7 (2)	9 (9)	19 (8)	3 (4)
H(4)	1031 (6)	42 (2)	-40 (2)	72 (13)	8 (3)	7 (2)	17 (10)	18 (9)	-2 (5)
H(5)	183 (4)	206 (2)	313 (2)	19 (9)	3 (2)*	5 (1)	5 (7)	-9 (6)	0 (3)*
H(6)	671 (5)	267 (2)	421 (2)	44 (11)	7 (2)	4 (2)	2 (8)	8 (7)	-3 (3)
H(7)	330 (6)	217 (2)	514 (2)	81 (15)	4 (3)	8 (2)	7 (11)	5 (9)	-3 (5)
H(8)	683 (4)	474 (2)	346 (2)	28 (10)	3 (2)	4 (2)	6 (7)	3 (6)	1 (3)
H(9)	531 (7)	503 (3)	544 (2)	71 (16)	11 (4)	13 (3)	1 (14)	-13 (12)	1 (6)
H(10)	135 (4)	518 (4)	383 (1)	19 (8)	5 (2)	3 (1)	-3 (6)	2 (5)	-4 (3)
H(11)	323 (5)	693 (2)	294 (2)	55 (13)	5 (3)	8 (2)	4 (10)	-6 (9)	-1 (5)
H(12)	395 (4)	603 (2)	193 (2)	37 (11)	9 (2)	7 (2)	11 (9)	11 (7)	6 (4)
H(13)	-24 (6)	726 (3)	170 (3)	62 (19)	11 (5)	18 (3)	9 (15)	5 (13)	-3 (7)

\* These parameters shifted by more than one e.s.d. during the last least-squares cycle in which the anisotropic temperature coefficients of the hydrogen atoms were adjusted. The largest shift was about 3.0 e.s.d. for  $b_{22}$  of H(1); the remaining shifts were less than 2.0 e.s.d.

Table 4. Hydrogen-bond distances and angles

Donor Atom	Acceptor Atom	Position of Acceptor Atom*	Distance		Angle
			D...A	H...A	
N(10)	N(7)	$-x+2, y-\frac{1}{2}, -z$	3.133 Å	2.32 Å	N(10)-H(3)-N(7) 167.1°
N(10)	O(5')	$-x+1, y-\frac{1}{2}, -z$	2.908	2.03	N(10)-H(4)-O(5') 169.0
O(2')	O(3')	$-x+1, y-\frac{1}{2}, -z+1$	2.749	1.90	O(2')-H(7)-O(3') 161.0
O(3')	N(3)	$-x+1, y+\frac{1}{2}, -z+1$	2.947	2.15	O(3')-H(9)-N(3) 149.3
O(5')	N(1)	$x-1, y+1, z$	2.774	1.92	O(5')-H(13)-N(1) 159.4

\*Relative to parameters of Table 2.

source of error. In addition, it is probable that the weights  $w$ , which were derived by adding together the weights of four or five individual observations (see *Experimental*), do not adequately reflect experimental errors other than counting statistics.

### Discussion

#### Hydrogen bonding and base stacking

Fig. 2 shows the structure viewed down the  $a$  axis. Each nucleoside is hydrogen bonded to six neighboring molecules; all nitrogen and oxygen atoms except N(9) and O(1') participate in the three-dimensional hydrogen bonding system. The hydrogen bond lengths and angles are given in Table 4. Additional close intermolecular contacts of 3.09 Å between C(2) and O(2') (at  $1-x, y-\frac{1}{2}, 1-z$ ) and 2.31 Å between H(1) and O(2') suggest a weak C-H...O bond; the angle H-C...O is 29°. No other intermolecular contact appears to be incompatible with accepted van der Waals radii.

The stacking of the purine rings in adenosine has already been described and compared with that in other crystalline derivatives of purine by Bugg (1971). The stacking involves molecules related to one another by translation along the  $a$  axis; the interplanar spacing between adjacent adenine moieties is 3.57 Å. The stacking interactions appear to involve primarily the amino group, which lies approximately above the midpoint of the C(5)-C(6) bond of the neighboring molecule.

#### Molecular geometry

The bond distances and angles involving non-hydrogen atoms are shown in Fig. 3 and those involving hydrogen atoms are given in Table 5. The esti-

mated standard deviations in distances involving the heavy atoms are about 0.003 Å and slightly less than 0.2° in the angles; for bonds involving hydrogen atoms, the uncertainties are approximately ten times larger.

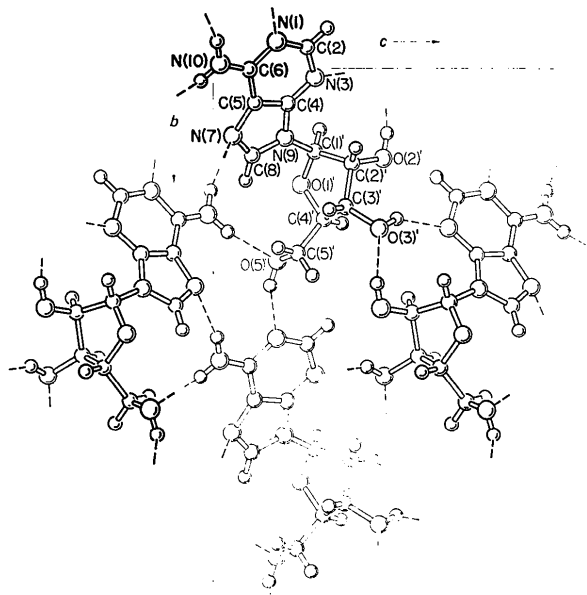


Fig. 2. The structure viewed down the  $a$  axis. Hydrogen bonds are indicated by dashed lines.

Table 5. Bond distances and angles involving hydrogen atoms

Standard deviations are about 0.03 Å and 2.0°

C(2)—H(1)	0.96 Å	C(3')—H(8)	0.94 Å
C(8)—H(2)	0.95	O(3')—H(9)	0.89
N(10)—H(3)	0.92	C(4')—H(10)	0.97
N(10)—H(4)	0.89	C(5')—H(11)	0.95
C(1')—H(5)	1.01	C(5')—H(12)	0.99
C(2')—H(6)	0.94	O(5')—H(13)	0.89
O(2')—H(7)	0.88		
H(1)—C(2)—N(1)	116.6°	H(8)—C(3')—C(2)	109.1°
H(1)—C(2)—N(3)	114.3	H(8)—C(3')—C(4')	110.2
H(2)—C(8)—N(7)	125.3	H(8)—C(3')—O(3')	105.8
H(2)—C(8)—N(9)	120.4	H(9)—O(3')—C(3')	104.5
H(3)—N(10)—H(4)	118.4	H(10)—C(4')—C(3')	108.0
H(3)—N(10)—C(6)	120.0	H(10)—C(4')—O(1')	108.7
H(4)—N(10)—C(6)	120.6	H(10)—C(4')—C(5')	111.6
H(5)—C(1')—N(9)	108.2	H(11)—C(5')—C(4')	108.9
H(5)—C(1')—C(2')	110.6	H(11)—C(5')—O(5')	107.4
H(5)—C(1')—O(1')	109.8	H(11)—C(5')—H(12)	110.1
H(6)—C(2')—C(1')	111.3	H(12)—C(5')—C(4')	111.6
H(6)—C(2')—C(3')	114.9	H(12)—C(5')—O(5')	109.4
H(6)—C(2')—O(2')	111.3	H(13)—O(5')—C(5')	109.2
H(7)—O(2')—C(2')	108.9		

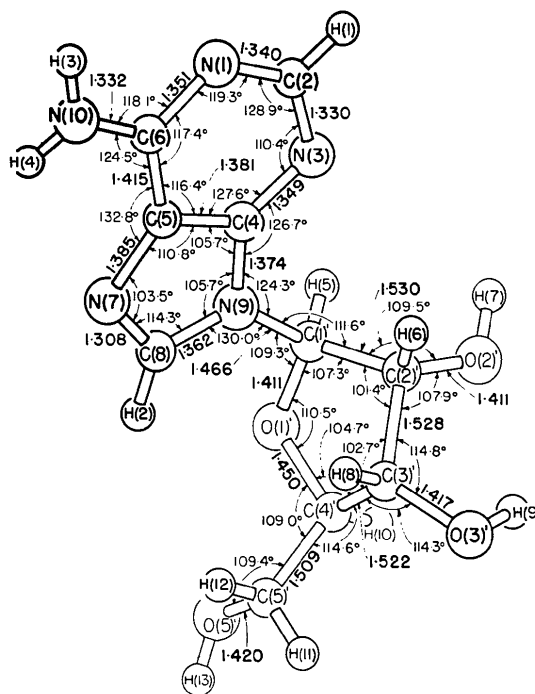


Fig. 3. Bond distances and angles involving the heavy atoms. Standard deviations are about 0.003 Å in the distances and 0.2° in the angles.

The reports of the crystal structures of ten compounds containing a neutral adenine ring substituted at the N(9) position have come to our attention: the pure compounds 9-methyladenine (Stewart & Jensen, 1964); deoxyadenosine (Watson, Sutor & Tollin, 1965); 3'-*O*-acetyladenosine (Rao, Sundaralingam & Arora, 1970);  $\alpha$ -D-2'-amino-2'-deoxyadenosine monohydrate (Rohrer & Sundaralingam, 1970); 1:1 complexes between 9-methyladenine and 1-methylthymine (Hoogsteen, 1963); 9-ethyladenine and 1-methyluracil (Mathews & Rich, 1964); 9-ethyladenine and 5-bromouracil (Katz, Tomita & Rich, 1966); 9-methyladenine and 1-methyl-5-bromouracil (Baklagine, Volkenstein & Kondrashev, 1966); 9-ethyladenine and 1-methyl-5-fluorouracil (Tomita, Katz & Rich, 1967); and a 2:1 complex between 8-bromo-9-ethyladenine and phenobarbital (Kim & Rich, 1968).

Within their uncertainties, the bond distances and angles in these other compounds are in satisfactory agreement with our values. In the five most accurate of the earlier analyses – the four uncomplexed compounds and the 9-methyladenine:1-methylthymine complex – the e.s.d.'s were in the range 0.005–0.010 Å and 0.3–0.5°; for these five determinations, the average values of the bond distances agree with ours within an r.m.s. deviation of 0.005 Å, and the angles within 0.6°. A common feature of all the structures is the relatively large difference between the bond angles N(1)–C(6)–N(10) and C(5)–C(6)–N(10) to the amino nitrogen atom.

Similarly, the bond lengths in the ribose moiety are in statistical agreement with those found in other nucleosides, including guanosine dihydrate and inosine dihydrate (Thewalt, Bugg & Marsh, 1970), adenosine-3'-phosphate (Sundaralingam, 1966), and 3'-*O*-acetyl-adenosine (Rao *et al.*, 1970) where the e.s.d.'s are in the range 0.004–0.006 Å and 0.2–0.3°. (The C(3')–O(3') distances are about 0.02 Å longer in the latter two

compounds, reflecting the effects of the attached phosphate and ester groups.) The bond angles, which are probably more sensitive to the conformations in the ribose rings, show larger differences; of particular note in the present structure are the relatively large values of the angles C(2')–C(3')–O(3') and C(4')–C(3')–O(3').

Average values for the C–H, N–H, and O–H distances are 0.96, 0.91, and 0.89 Å. These are typical for X-ray diffraction investigations, being somewhat shorter than the standard values for the internuclear separations.

A stereoscopic view of the molecule is shown in Fig. 4. Deviations from least-squares planes through the purine and the ribose moieties are listed in Table 6. It is interesting to note that the amino hydrogen atoms H(3) and H(4) are displaced from the plane of the adenine rings toward the hydrogen-bond acceptor atoms N(7) and O(5'). Similar, but statistically more significant, out-of-plane displacements of hydrogen atoms, presumably dictated by hydrogen bonding, have been noted in crystals of ethylenebis(biguanidine)-nickel(II) dichloride monohydrate (Holian & Marsh, 1970). On the other hand, atom H(1), which is involved in the intermolecular C(2)–H(1)···O(2') interaction, is displaced from the plane in a direction opposite from O(2').

A view of the molecule along the glycosidic bond C(1')–N(9) is shown in Fig. 5. Torsion angles in adenosine and in four related compounds are given in Table 7. The furanose ring is puckered with C(3') *endo*. The small value of the torsion angle  $\varphi$ (CN) about the glycosidic bond, 9.9°, is characteristic of purine nucleosides and nucleotides with 3' pucker of the ribose ring. Further discussions of torsion angles in related compounds are given by Sundaralingam (1969), among others.

The absolute configuration of the molecule has been

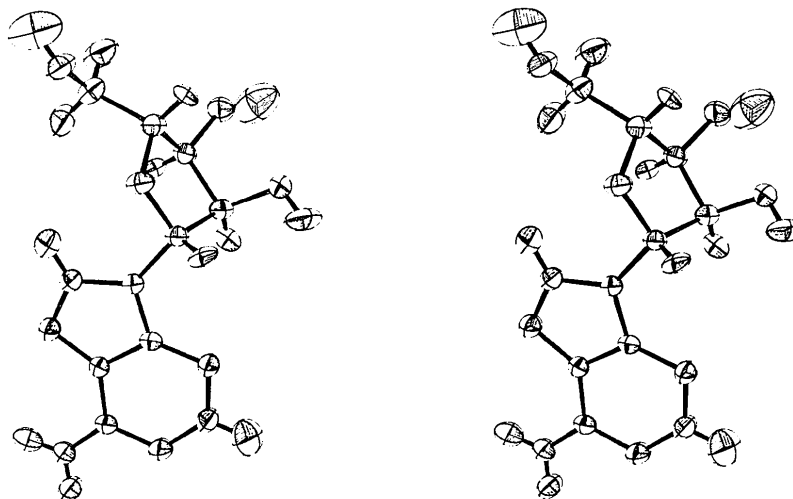


Fig. 4. A stereoscopic view of the molecule (Johnson, 1965).

confirmed by the use of anomalous scattering of Cu  $K\alpha$  radiation by oxygen and nitrogen atoms. Other reports in which the anomalous scattering by oxygen has been successfully used for this purpose are given by Hope & de la Camp (1969), Thiessen & Hope (1970), and Moncrief & Sims (1969).

#### Temperature ellipsoids

Ellipsoids representing the anisotropic temperature parameters are shown in Fig. 4. Thermal motions are reasonably small and isotropic for the heavy atoms, and in most instances the motions of the hydrogen atoms satisfactorily reflect those of the heavy atoms to which they are attached. Some of the hydrogen atoms have relatively large apparent motions parallel to the covalent bond, perhaps representing the delocalization of the bonding electron; we emphasize again, however,

Table 6. Deviations of the atoms from the least-squares planes

The direction cosines of the planes,  $q_i$ , are relative to the crystallographic axes;  $D$  is the origin-to-plane distance.

Ribose		Adenine	
C(1')	-0.039* Å	N(1)	0.021* Å
O(1')	0.041*	C(2)	-0.003*
C(2')	0.023*	N(3)	0.015*
C(3')	-0.554	C(4)	0.006*
C(4')	-0.025*	C(5)	0.019*
O(2')	1.371	C(6)	0.002*
O(3')	-0.240	N(7)	0.008*
C(5')	-0.897	C(8)	-0.016*
O(5')	-0.341	N(9)	-0.011*
N(9)	-1.290	N(10)	0.009
		C(1')	0.000
		H(1)	0.09
		H(2)	-0.02
		H(3)	-0.10
		H(4)	-0.04
		$q_a = 0.7395$	
		$q_b = 0.4471$	
		$q_c = 0.3772$	
		$D = 3.773$	
$q_a = -0.8156$			
$q_b = -0.1853$			
$q_c = 0.6727$			
$D = 0.654$ Å			

\* Atoms included in the calculation of the least-squares plane.

Table 7. Torsion angles in adenosine and related compounds

Angle	(a)	(b)	(c)	(d)	(e)
C(8)—N(9)—C(1')—O(1')	9.9°	12.4°	-139.9°	25.7°	3.8°
C(2')—C(1')—O(1')—C(4')	7.4	5.0	-26.4	4.8	-3.4
C(3')—C(2')—C(1')—O(1')	-27.3	-28.7	37.1	-29.8	-19.1
C(4')—C(3')—C(2')—C(1')	35.7	38.5	-33.2	42.3	34.6
O(1')—C(4')—C(3')—C(2')	-32.5	-36.9	18.8	-40.0	-37.6
C(1')—O(1')—C(4')—C(3')	16.0	20.7	4.8	22.8	25.6
C(3')—C(4')—C(5')—O(5')	177.0	40.9	58.0	40.0	171.2

(a) Adenosine; this investigation.

(b) Adenosine-5-bromouridine (Haschemeyer & Sobell, 1963; Sundaralingam, 1969).

(c) 3'-O-Acetyladenosine (Rao, Sundaralingam & Arora, 1970).

(d) Adenosine-5'-phosphate (Kraut & Jensen, 1963; Sundaralingam, 1969).

(e) Adenosine-3'-phosphate dihydrate (Sundaralingam, 1966; 1969).

that the temperature parameters of the hydrogen atoms have not been determined with high precision.

#### References

- BAKLAGINE, Y. G., VOLKENSTEIN, M. V. & KONDRASHEV, Y. D. (1966). *Zh. Struk. Khim.* **7**, 399.
- BUGG, C. E. (1971). *Proceedings, Fourth Jerusalem Symposium on the Purines: Theory and Experiment*. Edited by B. PULLMAN. New York: Academic Press.
- CROMER, D. T. & LIBERMAN, D. (1970). *J. Chem. Phys.* **53**, 1891.
- DUCHAMP, D. J. (1964). Amer. Cryst. Assoc. Meeting, Bozeman, Montana, paper B-14, p. 29.
- HASCHEMEYER, A. E. V. & SOBELL, H. M. (1965). *Acta Cryst.* **18**, 525.
- HOLIAN, B. L. & MARSH, R. E. (1970). *Acta Cryst.* **B26**, 1049.
- HOOGSTEEN, K. (1963). *Acta Cryst.* **16**, 907.
- HOPE, H. (1969). Personal communication.
- HOPE, H. & DE LA CAMP, U. (1969). *Nature*, **221**, 54.
- International Tables for X-ray Crystallography* (1962). Vol. III, p. 202. Birmingham: Kynoch Press.
- JOHNSON, C. K. (1965). *ORTEP*. ORNL-3794, Oak Ridge National Laboratory, Oak Ridge, Tennessee, U.S.A.
- KATZ, L., TOMITA, K. & RICH, A. (1966). *Acta Cryst.* **21**, 754.
- KIM, S. H. & RICH, A. (1968). *Proc. Natl. Acad. Sci. U.S.* **60**, 402.
- KRAUT, J. & JENSEN, L. H. (1963). *Acta Cryst.* **16**, 79.
- LARSON, A. C. (1967). *Acta Cryst.* **23**, 664.
- MATHEWS, F. S. & RICH, A. (1964). *J. Mol. Biol.* **8**, 89.
- MONCRIEF, J. W. & SIMS, S. P. (1969). *Acta Cryst.* **A25**, supplement, page S74.

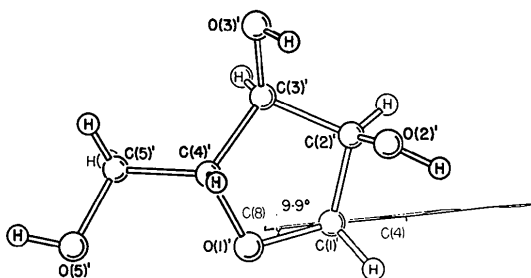


Fig. 5. The molecule as viewed along the glycosidic bond C(1')—N(9).

- RAO, S. T., SUNDARALINGAM, M. & ARORA, S. K. (1970). *J. Amer. Chem. Soc.* **92**, 4963.
- ROHRER, D. C. & SUNDARALINGAM, M. (1970). *J. Amer. Chem. Soc.* **92**, 4956.
- STEWART, R. F. & JENSEN, L. H. (1964). *J. Chem. Phys.* **40**, 2071.
- STEWART, R. F., DAVIDSON, E. R. & SIMPSON, W. T. (1965). *J. Chem. Phys.* **42**, 3175.
- SUNDARALINGAM, M. (1969). *Biopolymers* **7**, 821.
- SUNDARALINGAM, M. (1966). *Acta Cryst.* **21**, 495.
- THEWALT, U., BUGG, C. E. & MARSH, R. E. (1970). *Acta Cryst.* **B26**, 1089.
- THIESSEN, W. E. & HOPE, H. (1970). *Acta Cryst.* **B26**, 554.
- TOMITA, K., KATZ, L. & RICH, A. (1967). *J. Mol. Biol.* **30**, 545.
- WATSON, D. G., SUTOR, D. J. & TOLLIN, P. (1965). *Acta Cryst.* **19**, 111.
- WILSON, A. J. C. (1942). *Nature, Lond.* **150**, 152.
- ZACHARIASEN, W. H. (1963). *Acta Cryst.* **16**, 1139.

*Acta Cryst.* (1972). **B28**, 1989

## The Crystal Structure of Diiodoacetylene; an Example of Pseudosymmetry

BY J. D. DUNITZ, HERMANN GEHRER AND DOYLE BRITTON\*

Laboratory for Organic Chemistry, Swiss Federal Institute of Technology, 8006 Zürich, Switzerland

(Received 18 January 1972)

Diiodoacetylene,  $C_2I_2$ , forms tetragonal crystals with  $a=15.68(3)$ ,  $c=4.30(1)$  Å, eight molecules in the unit cell, space group  $P4_2/n$ . The structure was determined from three-dimensional diffractometer data, and the refinement carried to a conventional  $R$  of 0.145. The problem was complicated by decomposition of the crystals, twinning and pseudosymmetry in the diffraction pattern. The linear molecules are in general positions, but each molecule points, at both ends, towards the middles of two adjacent molecules to give four  $C \cdots I$  distances of about 3.4 Å, which we interpret as donor( $C \equiv C$ )–acceptor( $I$ ) interactions. The intramolecular  $I \cdots I$  distance is 5.16 Å.

### Introduction

Several years ago Professor Jerry Donohue pointed out to one of us an apparently anomalous situation in the structure of crystalline diiodoacetylene,  $C_2I_2$ . Flannery & Nixon (Flannery, 1966) had made X-ray diffraction photographs that showed Laue symmetry  $4/mmm$  and systematic extinctions corresponding to an  $n$  glide plane perpendicular to the  $c$  axis. This evidence would require the space group to be  $P4/nmm$ , although there was also an indication of a  $4_2$  axis parallel to  $c$ . The cell dimensions,  $a=b=15.65$ ,  $c=4.29$  Å, and measured density suggested eight molecules in the unit cell. No hypothetical structure consistent with the  $P4/nmm$  space group with  $Z=8$  could be found, and an examination of the  $hk0$  reflexions (Patterson projection along the short  $c$  axis) did not provide a solution. The possibility of twinning was raised but not explored further. Additional complications are caused by gradual decomposition of the compound in the X-ray beam.

We have studied the problem further and report here our results.

### Experimental

A sample of diiodoacetylene was supplied to us by Dr E. Kloster-Jensen (Kloster-Jensen, 1969). Crystals

suitable for X-ray diffraction studies were prepared by sublimation at 50°C under reduced pressure; since the crystals are light-sensitive it was necessary to work in a darkened room. Precession photographs (Mo  $K\alpha$  radiation,  $\lambda=0.7107$  Å) showed the crystals to be tetragonal with  $a=15.68(3)$  and  $c=4.30(1)$  Å, in agreement with Flannery's values of 15.65 and 4.29 Å. Careful scrutiny of the photographs shows that the  $P4/mmm$  symmetry is only approximate; about 20% of the reflexion-pairs  $hkl$  and  $khl$  show small but significant intensity differences that cannot be explained away as absorption effects. The intensity differences vary from crystal to crystal but they affect the same reflexion pairs. The true Laue symmetry is only  $P4/m$ . Systematic extinctions:  $0k0$ ,  $k=2n+1$ ;  $hk0$ ,  $h+k=2n+1$  show the space group to be  $P4_2/n$ . (The following discussion is based on the second arrangement of  $P4_2/n$  given in *International Tables for X-ray Crystallography* (1959) with the origin at  $\bar{1}$ .) The density of 3.43(3) g.cm<sup>-3</sup> measured by flotation in thallos formate–thallous malonate solution agrees with the value of 3.49(1) g.cm<sup>-3</sup> calculated for  $Z=8$ .

Intensity data were collected using the Hilger and Watts linear diffractometer (Mo  $K\alpha$  radiation, balanced filters, crystal mounted with  $c$  as the rotation axis, two cycles per reflexion). At room temperature, at which all the measurements were made, the crystals decomposed in the X-ray beam at such a rate that it was necessary to use three different crystals to collect the data; one for  $hk0$ , one for  $hk1$  and  $hk2$ , and one for  $hk3$  and  $hk4$ . All the crystals were prisms elongated

\*Permanent address: Chemistry Department, University of Minnesota, Minneapolis 55455, U.S.A.